

# ENERGETIC PARTICLES AND BURNING PLASMA PHYSICS

Reported by

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and Integrated Simulation of Burning Plasmas  
& 21st Century COE Workshop on Plasma Theory  
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## Reporting work by:

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### REFERENCES

- A. Ödblom, B. N. Breizman, S. E. Sharapov, T. C. Hender, and V. P. Pastukhov, *Nonlinear magnetohydrodynamical effects in precessional fishbone oscillations*, Phys. Plasmas **9**, 155 (2002)
- Y. Todo, H. L. Berk, and B. N. Breizman, *Simulation of intermittent beam ion loss in a Tokamak Fusion Test Reactor experiment*, Phys. Plasmas **10**, 2888 (2003)
- B. N. Breizman, H. L. Berk, and M. S. Pekker, *Theory of Alfvén eigenmodes in shear reversed plasmas*, Phys. Plasmas **10**, 3649 (2003)
- B. N. Breizman, *Selected issues in energetic particle theory*, 8th IAEA Technical Committee Meeting on Energetic Particles in Magnetic Confinement Systems (October 2003, San Diego)

# INTRODUCTION

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- **Alpha particles in a burning plasma can have an affect on:**
  - Stability: fishbone, TAE
  - Transport: particle loss, wall heat loading
  - Heating and current drive
  - Edge physics: transport barrier
  - Burn dynamics: He ash, thermal stability
- **Hence the need for integrated simulation, coupled with theory-based modeling, of alpha particle physics in a burning plasma**

# ISSUES FOR ALPHA THEORY/SIMULATION

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- **Predictive capabilities:**
  - Limitations of linear theory/simulations
  - Near-threshold nonlinear regimes (soft/hard)
- **Alfvén Cascade modes (example of soft nonlinear response)**
  - Interpretation in reversed-shear scenarios
  - Mysteries that remain to be explained
- **Fishbone modes (example of hard nonlinear response)**
  - Interplay of fluid (continuum) and kinetic (wave-particle) resonances
- **Convective and diffusive scenarios for global transport**
  - Transport barriers for energetic/alpha particles

# PREDICTIVE CAPABILITIES

# LIMITATIONS OF LINEAR THEORY/SIMULATION

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## “ROBUST”

- **Mode frequencies**
  - Robust (for perturbative modes)
  - Easy to measure experimentally
- **Mode structure**
  - Robust, experimentally measurable, and useful in nonlinear theory (for perturbative modes)

## “LESS ROBUST”

- **Growth rate**
  - Energetic particle drive can be reliably calculated
  - Useful as parameter in near-threshold weakly nonlinear theory
  - However, full nonlinear effects may cause it to change significantly
- **Damping rates**
  - Except for collisional damping, the damping rates from the background plasma (e.g., ion Landau damping, radiative damping, continuum damping) are exponentially sensitive to parameter/profile details
  - Exp'tally measurable (for reverse determination of parameters)

# UTILITY OF LINEAR THEORY/SIMULATION

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- **Study the phase space for fast particles in the presence of linearly unstable modes, to find resonance overlap conditions for global diffusion**
  - Note: Possible partial depletion of phase space (“holes” in distribution)
- **Study the sensitivity of instability boundaries to the variation of plasma parameters**
  - Difficult to do globally; instead analyze selected modes for trends
- **Study whether unstable modes can be suppressed over a radial interval that is sufficiently broad to create a transport barrier for energetic particles**
- **Assess edge effects (e.g., boundary conditions for the modes, particle orbits, etc.)**

# NEAR-THRESHOLD NONLINEAR REGIMES

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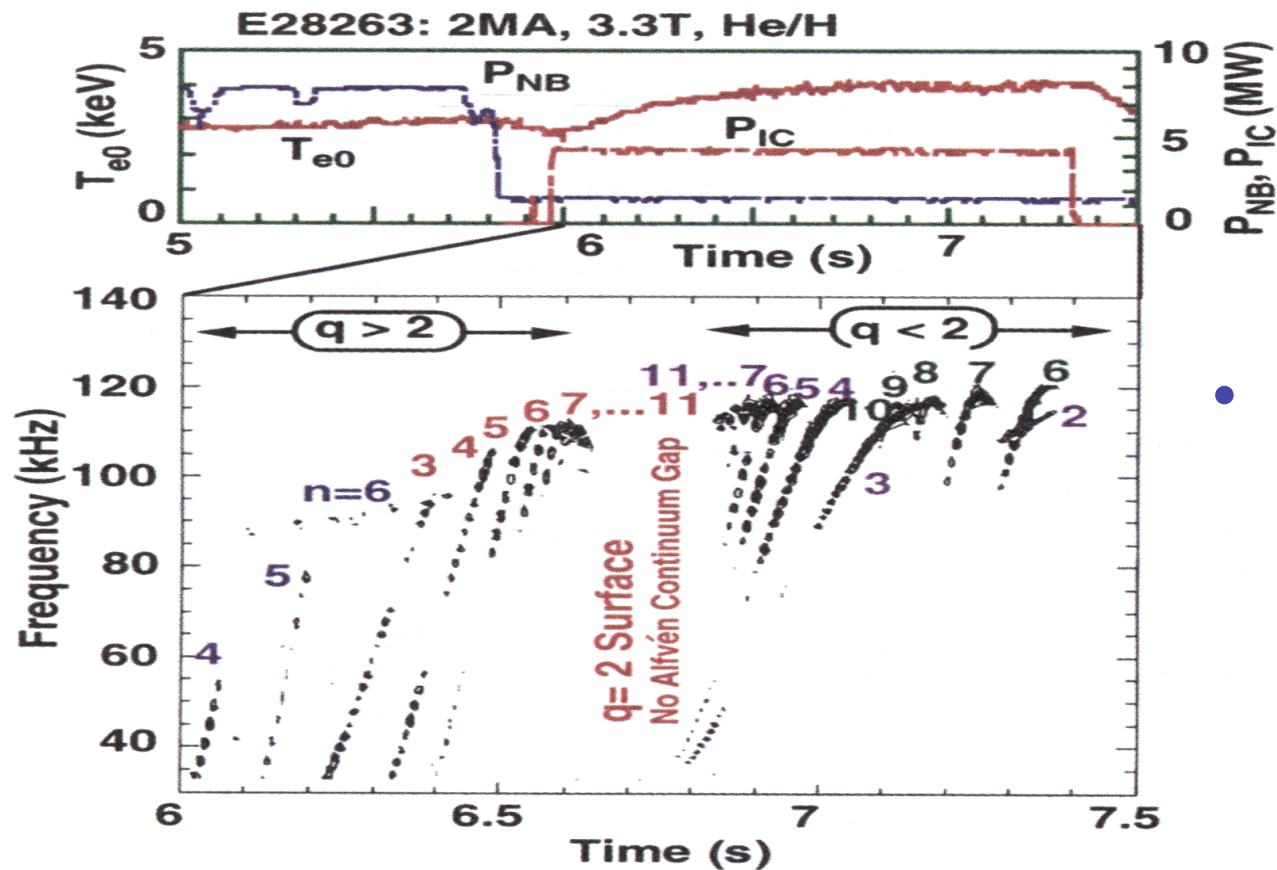
- **Why study the nonlinear response near the threshold?**
  - Typically, macroscopic plasma parameters evolve slowly compared to the instability growth time scale
  - Perturbation technique is adequate near the instability threshold
- **Single-mode case:**
  - Identification of the soft and hard nonlinear regimes is crucial to determining whether an unstable system will remain at marginal stability
  - Bifurcations at single-mode saturation can be analyzed
  - The formation of long-lived coherent nonlinear structure is possible
- **Multi-mode case:**
  - Multi-mode scenarios with marginal stability (and possibly transport barriers) are interesting
  - Resonance overlap can indicate hard nonlinear regime



# ALFVEN CASCADES

(example of soft nonlinear response)

# OBSERVATION OF ALFVEN CASCADES

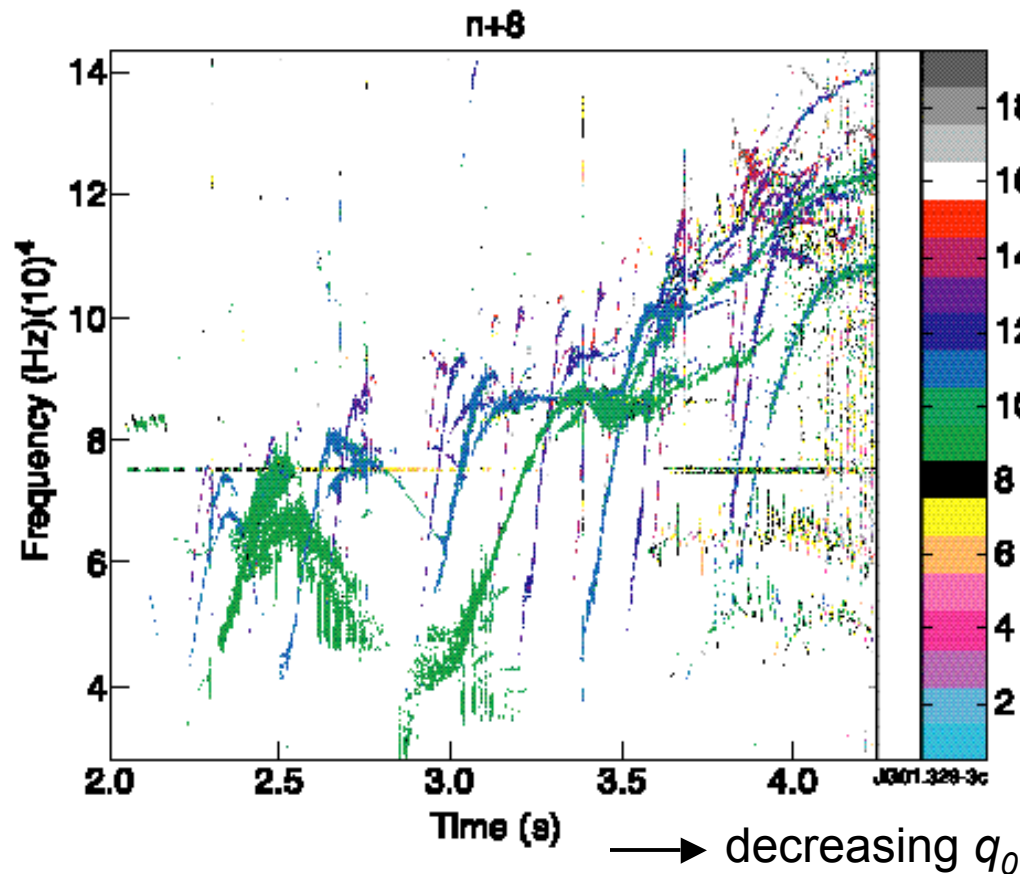


- First seen in JT-60U
  - Fukuyama simulation of “Reversed Shear Induced Alfvén Eigenmodes”

H. Kimura et al., Nuclear Fusion **38**, 1303 (1998)

→ decreasing q

# FEATURES OF ALFVEN CASCADES

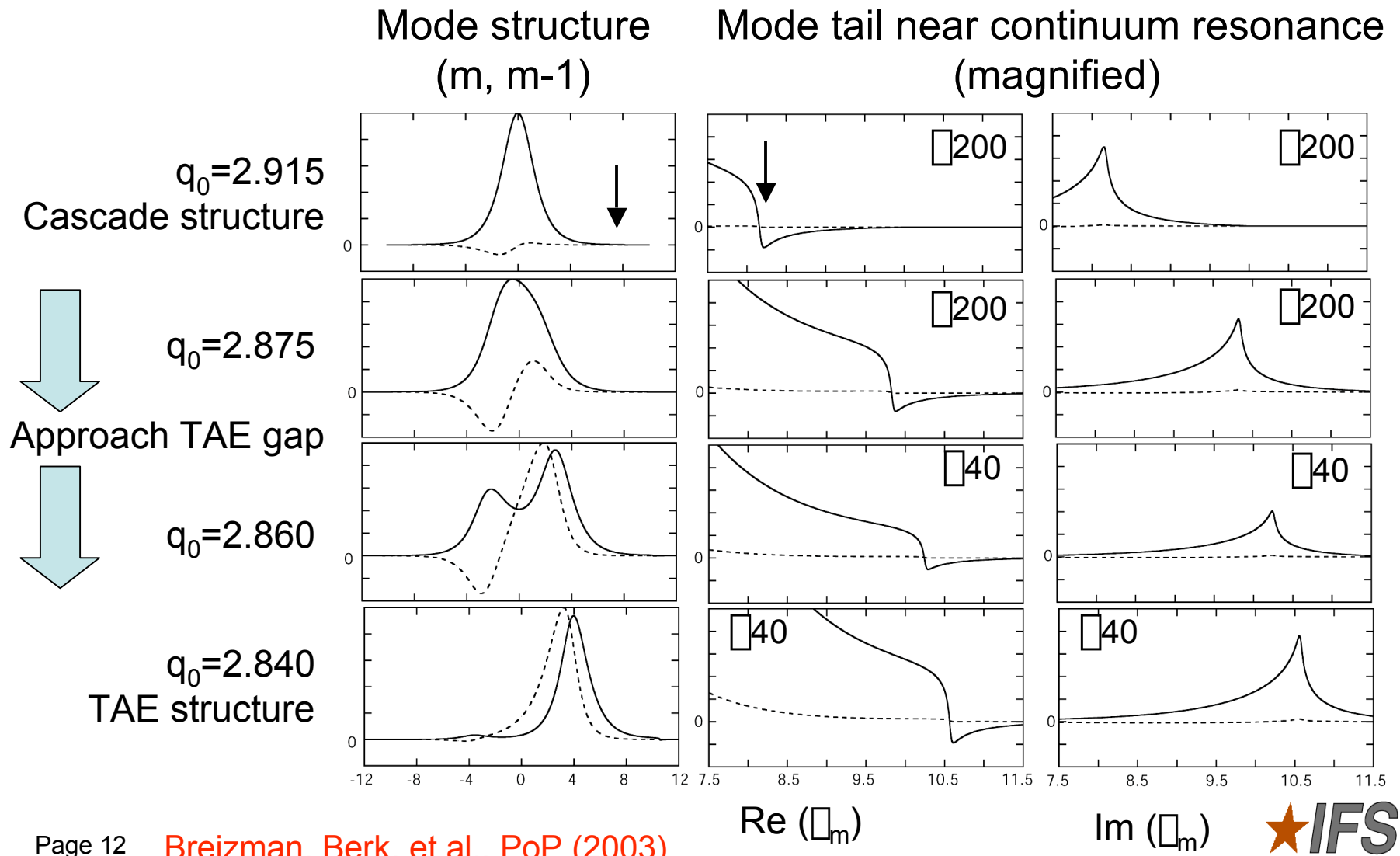


S. Sharapov and IFS/JET teams,  
Phys Lett A **289**, 127 (2001)

## • Characteristics:

- Occur in reversed shear plasmas at zero-shear location  $q(r) = q_0$
- Frequency initially below TAE gap and sweeps upward (usually never downward) as safety factor decreases in time
- Frequency signals are quasi-periodic for a large number of  $n$ -values (typically,  $n=1-6$ )
- Higher- $n$  modes recur more frequently and have more rapid frequency sweeping ( $\partial f/\partial t \propto n$ )
- Sometimes the frequency merges into the TAE gap
- Modes are suppressed at low frequencies

# TRANSITION OF ALFVEN CASCADES TO TAEs



# CONTINUUM DAMPING FOR CASCADE MODES

- Cascade eigenmode equation (with  $\propto x^4$  continuum damping terms):

$$\frac{d}{dx} \left[ (S + i\gamma) + x^2 \propto x^4 \right] \frac{d}{dx} \psi_m \propto \left[ (S + i\gamma) + x^2 \propto \frac{Q}{2(m - nq_0)} \propto x^4 \right] \psi_m = 0$$

$$(S + i\gamma) \equiv \frac{\propto (\propto + i\gamma) \propto 4(m - nq_0)^2}{4(m - nq_0)} \frac{mq_0}{r_0^2 q_0 \propto} \quad \propto \equiv \frac{1}{4(m - nq_0)} \frac{r_0^2 q_0 \propto}{mq_0} \ll 1$$

Energetic particle contribution:  $Q \equiv \propto \frac{4\propto e}{cB} \frac{q_0^2}{q_0} V_A \frac{R}{r_0} \frac{d}{dr} \langle n_{\text{fast ion}} \rangle$

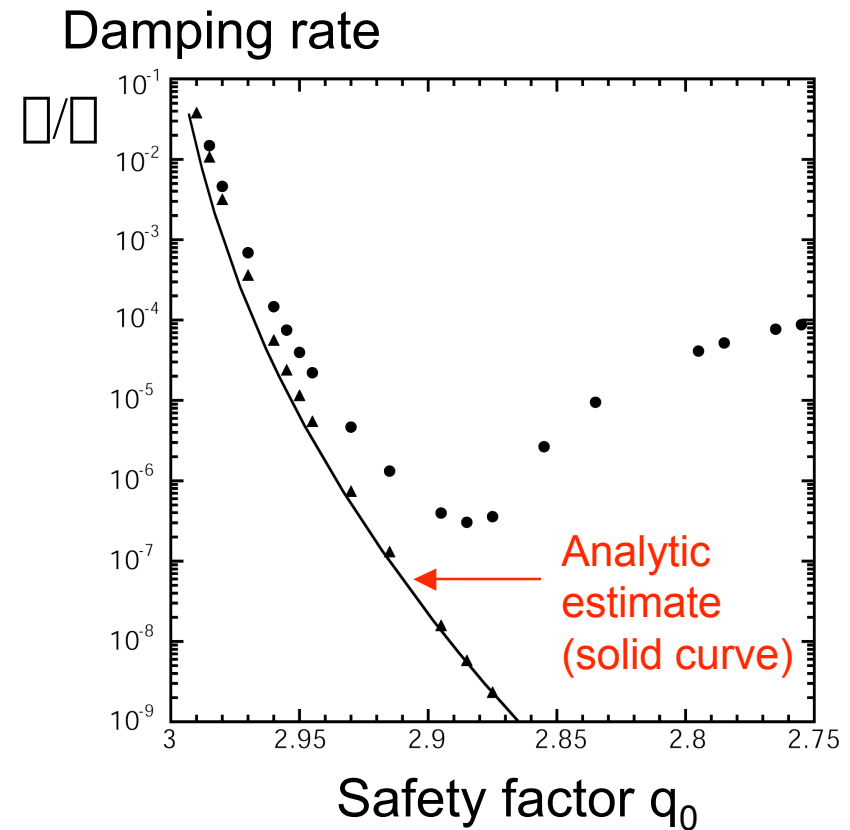
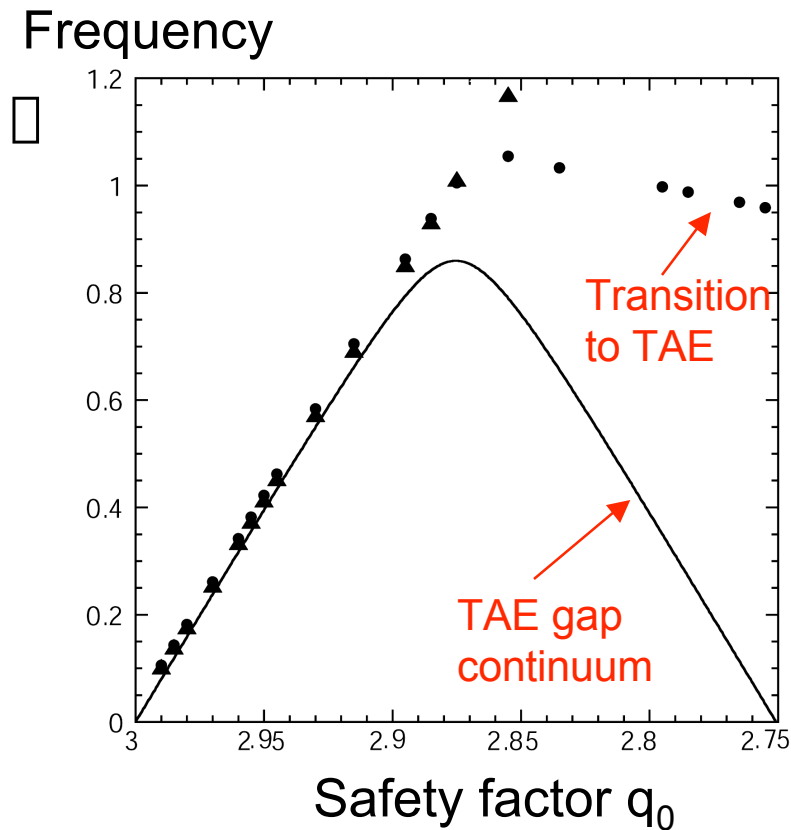
- Damping rate grows exponentially as  $q_0$  approaches rational surface:

$$\text{Im } \propto \propto \exp\left(\propto \frac{2}{\sqrt{\propto}}\right)$$

- Low frequency limit:  $\propto_{\min} = 2Q \frac{r_0^2 q_0 \propto}{mq_0} \rightarrow \propto_{\min} = \propto \frac{\propto_{ci}}{m} \frac{r_0}{n} \frac{d}{dr} \langle n_{\text{fast ion}} \rangle$

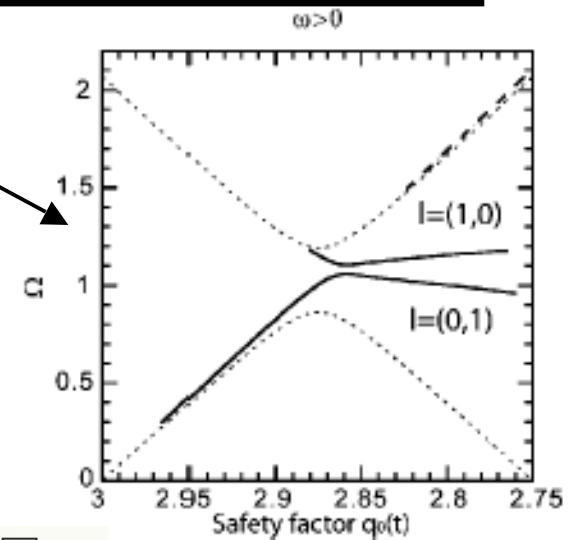
# CASCADE FREQUENCY AND DAMPING RATE

- ▲ = single poloidal harmonic ( $m = 12$ )
- = two coupled poloidal harmonics ( $m = 11, m = 12$ )

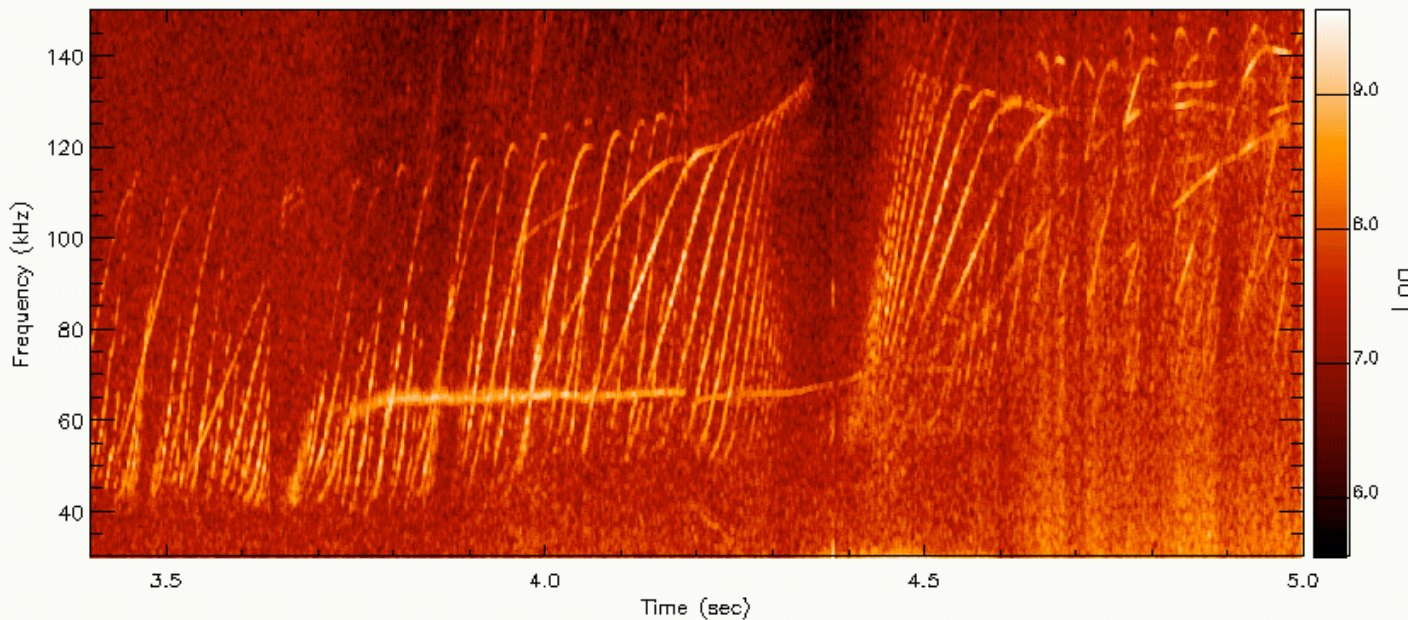


# PARTIALLY EXPLAINED FEATURES

- **Transition from Alfvén Cascades to TAEs**
  - But why not continue after turn over?
- **Suppression of Cascades at low frequency:**
  - Continuum damping for moderate  $m$  (radiative damping at high  $m$ )
  - Coupling to acoustic modes & ion Landau damping?
  - Low-freq cutoff could provide a diagnostic tool



Breizman, Berk, et al.  
Phys Plasmas (2003)

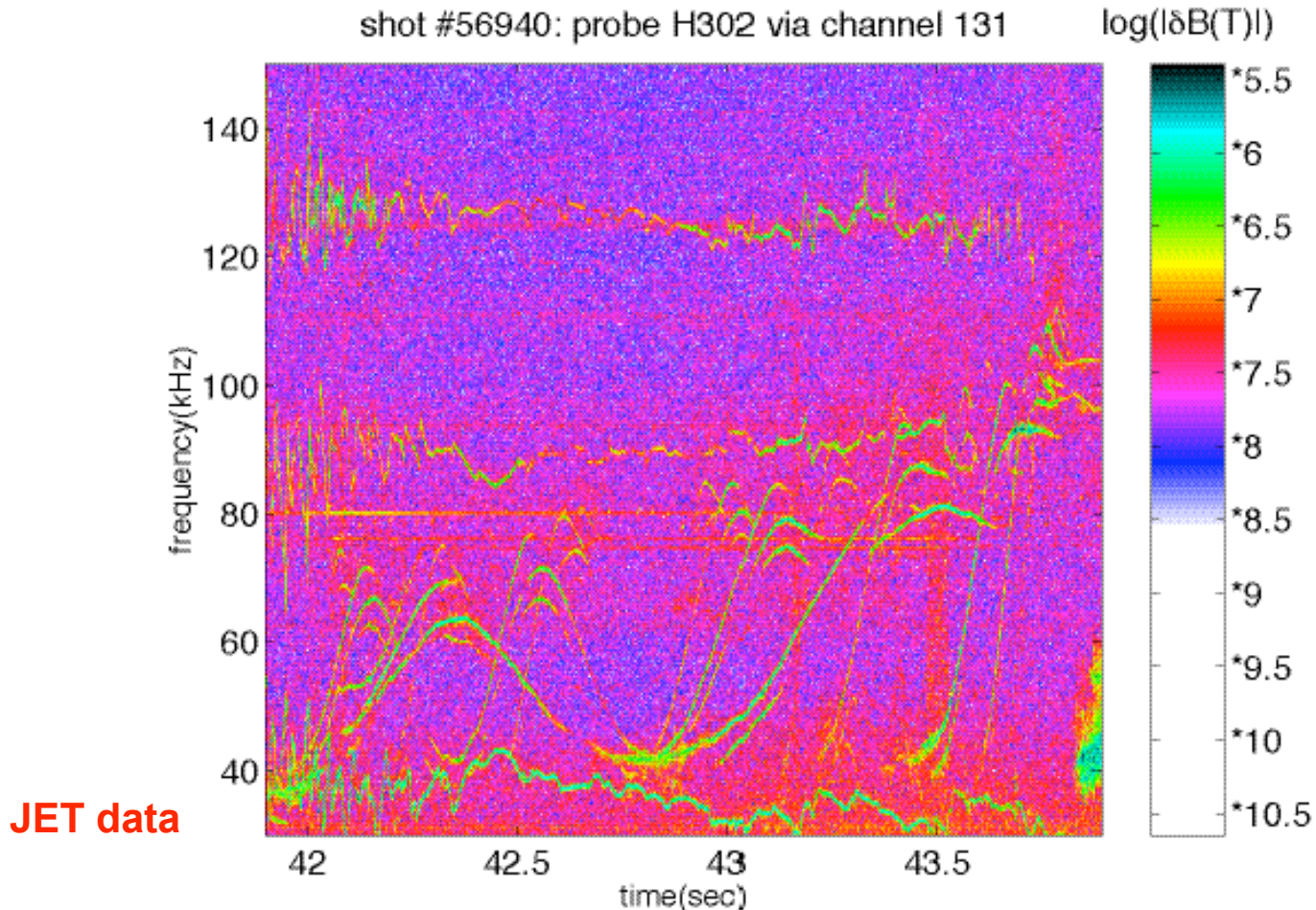


JET data



# UNEXPLAINED CASCADE FEATURES

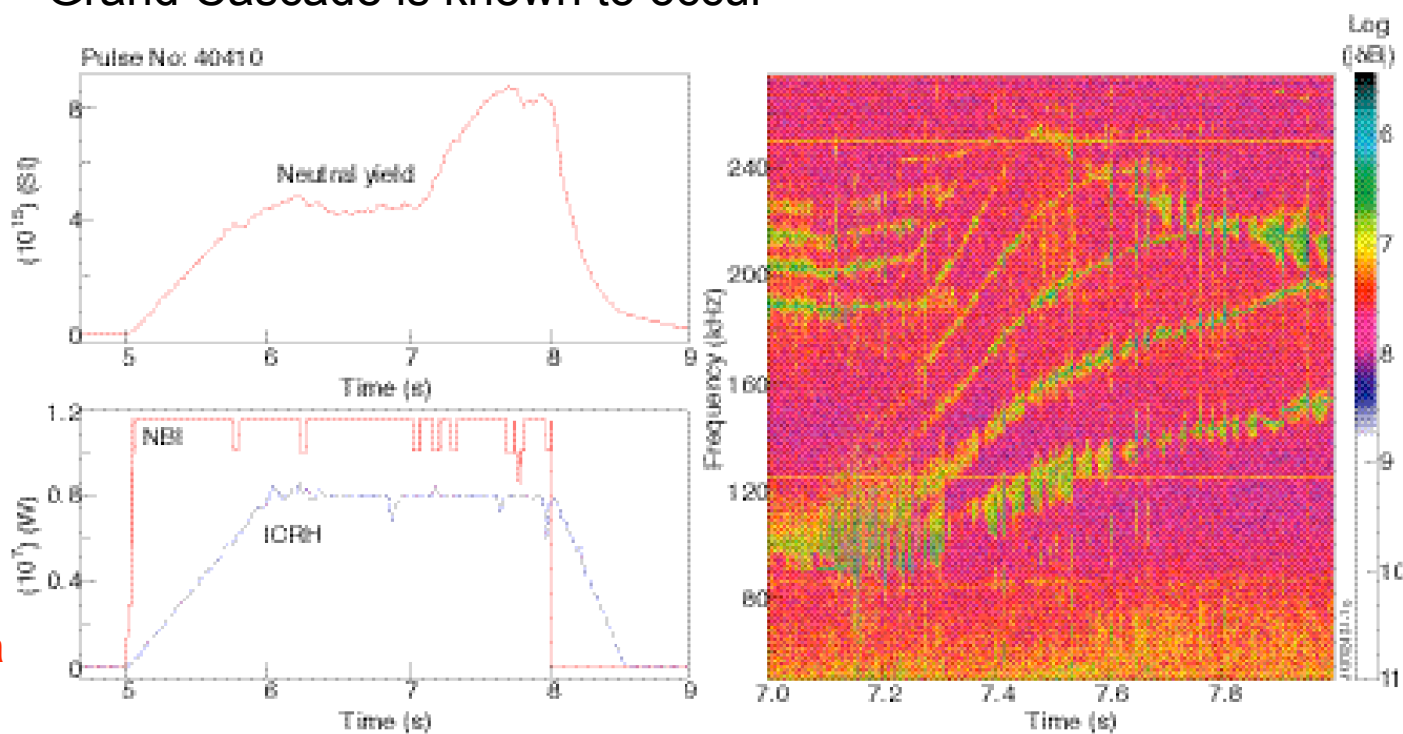
- Frequency “rolling”
- Mode enhancement at low frequencies





# ALFVEN CASCADES AS ITB DIAGNOSTIC

- **ITB triggering event**
  - “Grand Cascade” (many simultaneous n-modes) occurrence is coincident with ITB formation (when  $q_{\min}$  passes through integer value)
  - Already being used on JET as a diagnostic to monitor  $q_{\min}$
  - Proposal to create ITB by application of main heating shortly before a Grand Cascade is known to occur



JET data

# FISHBONES

(example of hard nonlinear response)

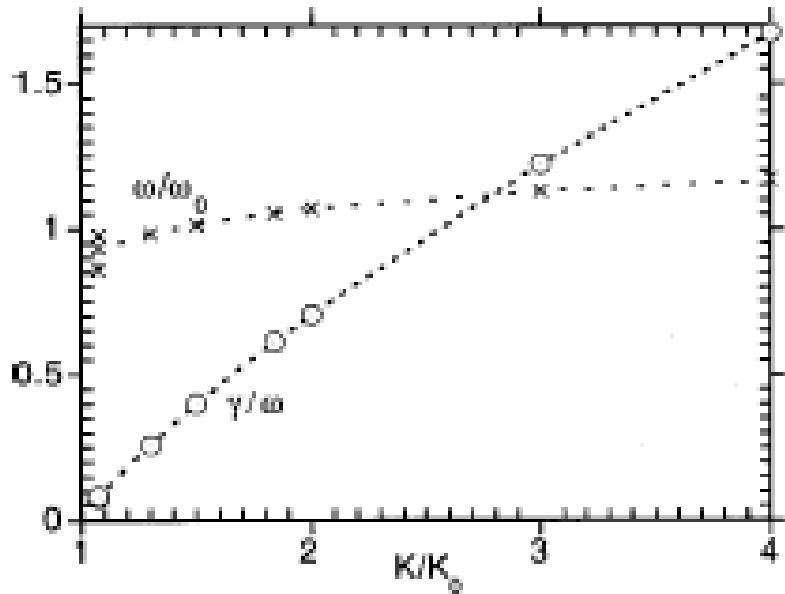
# FISHBONE ONSET

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- **Linear responses from kinetic (wave-particle) and fluid (continuum) resonances are in balance at instability threshold; however, their nonlinear responses differ significantly.**
- **Questions:**
  - Which resonance produces the dominant nonlinear response?
  - Is this resonance stabilizing or destabilizing?
- **Approach:**
  - Analyze the nonlinear regime near the instability threshold
  - Perform hybrid kinetic-MHD simulations to study strong nonlinearity

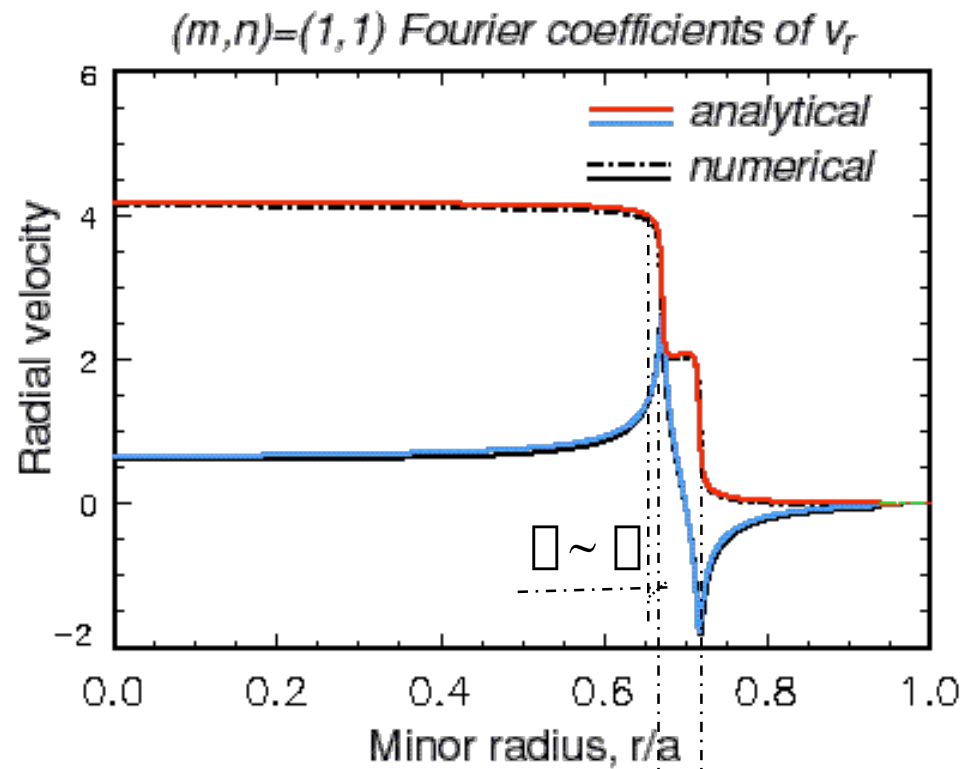
# LINEAR NEAR-THRESHOLD MODE

## Frequency $\omega$ and growth rate $\gamma$



Energetic ion drive

## Double resonance layer at $\omega = \pm \omega(r)$



$$\omega = r \omega_A S^{\frac{1}{2}}$$

A. Odblom et al., Phys Plasmas (2002)

# EARLY NONLINEAR DYNAMICS

- **Weak MHD nonlinearity of the  $q = 1$  surface destabilizes fishbone perturbations**
  - Near threshold, fluid nonlinearity dominates over kinetic nonlinearity
  - As mode grows,  $q$  profile is flattened locally (near  $q=1$ )  $\square$  continuum damping reduced  $\square$  **explosive growth** is triggered

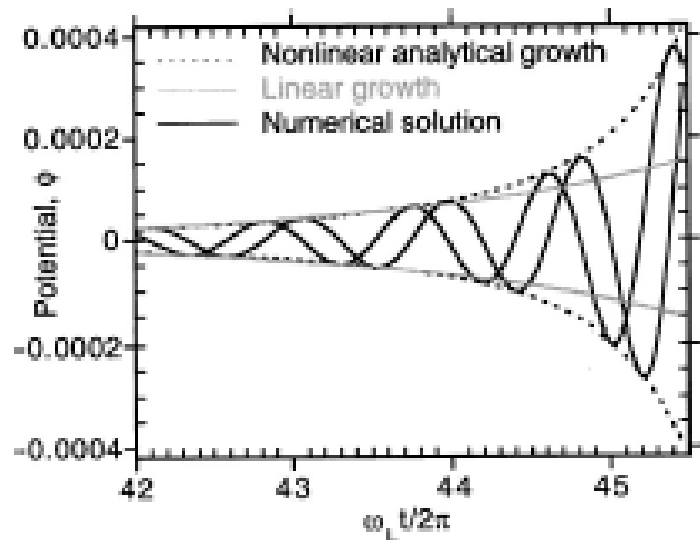


FIG. 4. The  $m=1$  amplitude evolution in the numerical solution (solid black), and as predicted by the explosive analytical envelope (dashed), in contrast to a linear growth (solid gray), for a mode near the instability threshold,  $\gamma/\omega \sim 8\%$ , with  $m=0-8$ ,  $\eta=0=\nu$ ,  $T/\tau_A=40$ .

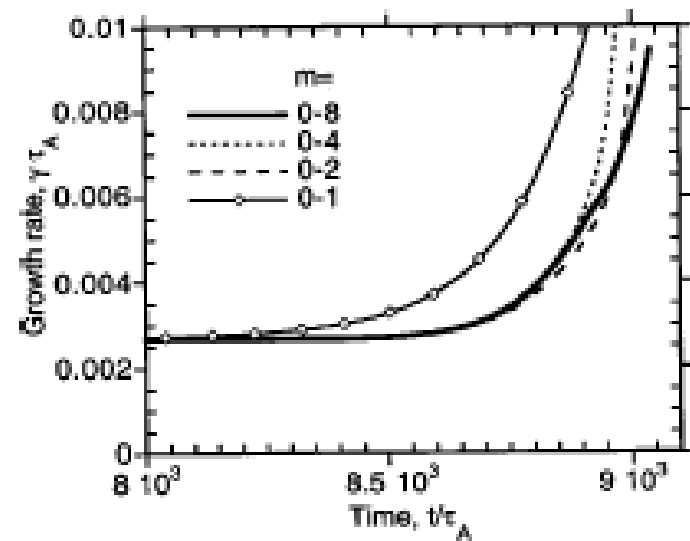


FIG. 5. The nonlinear growth rate evolution of the  $m=1$  radial velocity amplitude,  $\gamma = \partial \ln|V_r| / \partial t$ , for a mode near the instability threshold,  $\gamma/\omega \sim 8\%$ , and its dependence on the numerical spectral resolution:  $m=0-8$  (thick solid),  $m=0-4$  (short dashed),  $m=0-2$  (long dashed),  $m=0-1$  (solid with markers). Here,  $T/\tau_A=40$ .

# UNEXPLAINED FISHBONE FEATURES

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- **Transition from explosive growth to slowly growing MHD structure (i.e., island near  $q=1$  surface)**
- **Modification of fast particle distribution**
- **Mode saturation and decay**
- **Quantitative simulation of frequency sweeping**
  - Frequency change during explosive phase suggests mode will slow down and saturate
- **Burst repetition rate in presence of injection**
  - Need to include sources/sinks/collisions

# CONVECTIVE AND DIFFUSIVE TRANSPORT

# CONVECTIVE TRANSPORT IN PHASE SPACE

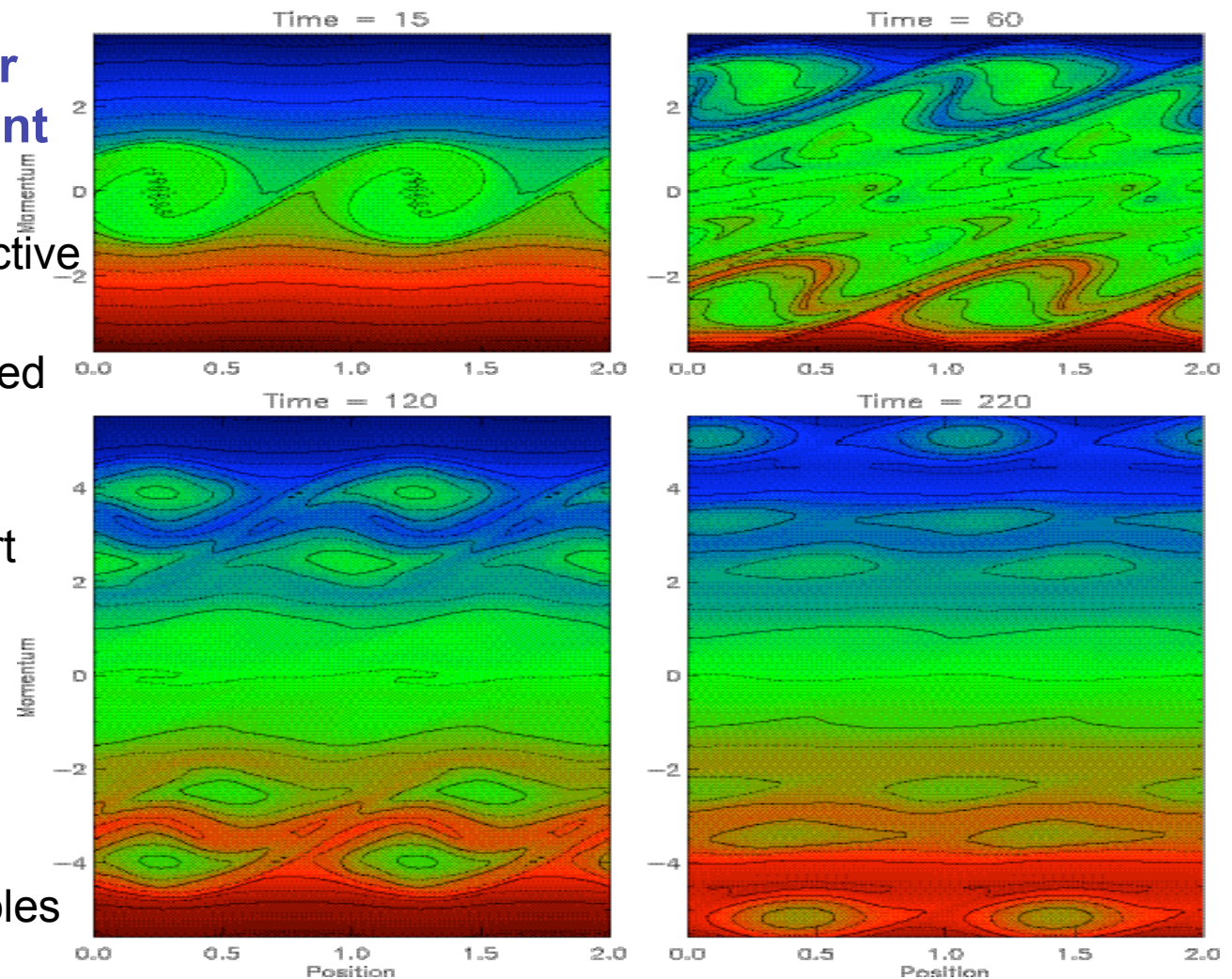
PHASE SPACE PLOTS OF PARTICLE DISTRIBUTION

- **Explosive behavior can lead to coherent structures**

- Can cause convective transport
- Single mode: limited extent
- Multiple modes: extended transport (“avalanche”)

- **Not simulated for fishbone**

- Indication from another example (single mode): “holes and clumps”



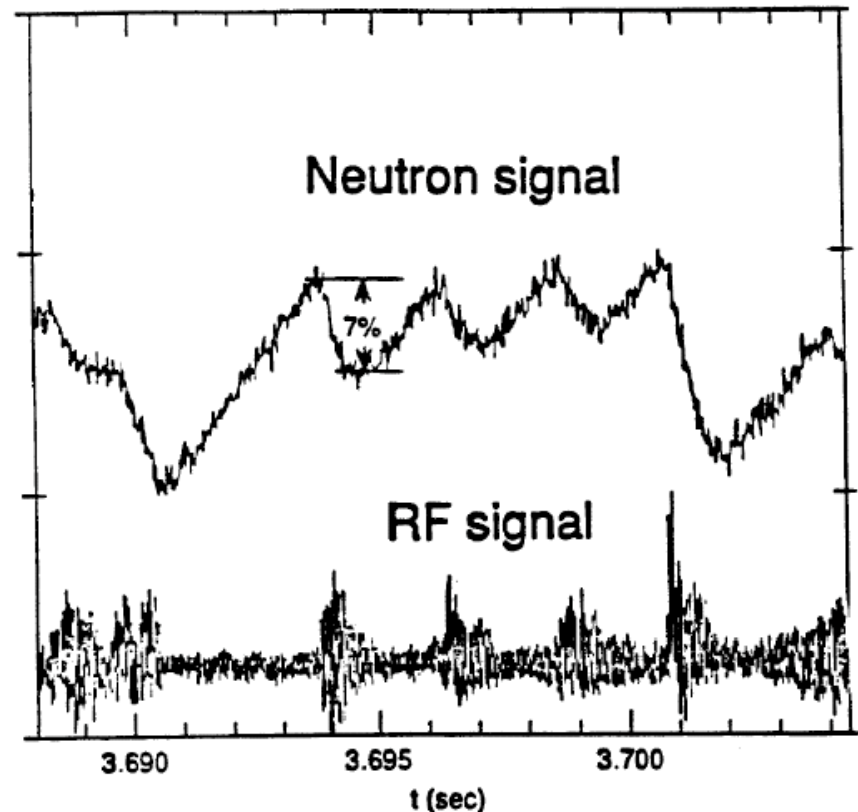
Petviashvili, Berk, Breizman (1998)





# INTERMITTENT LOSSES OF FAST IONS

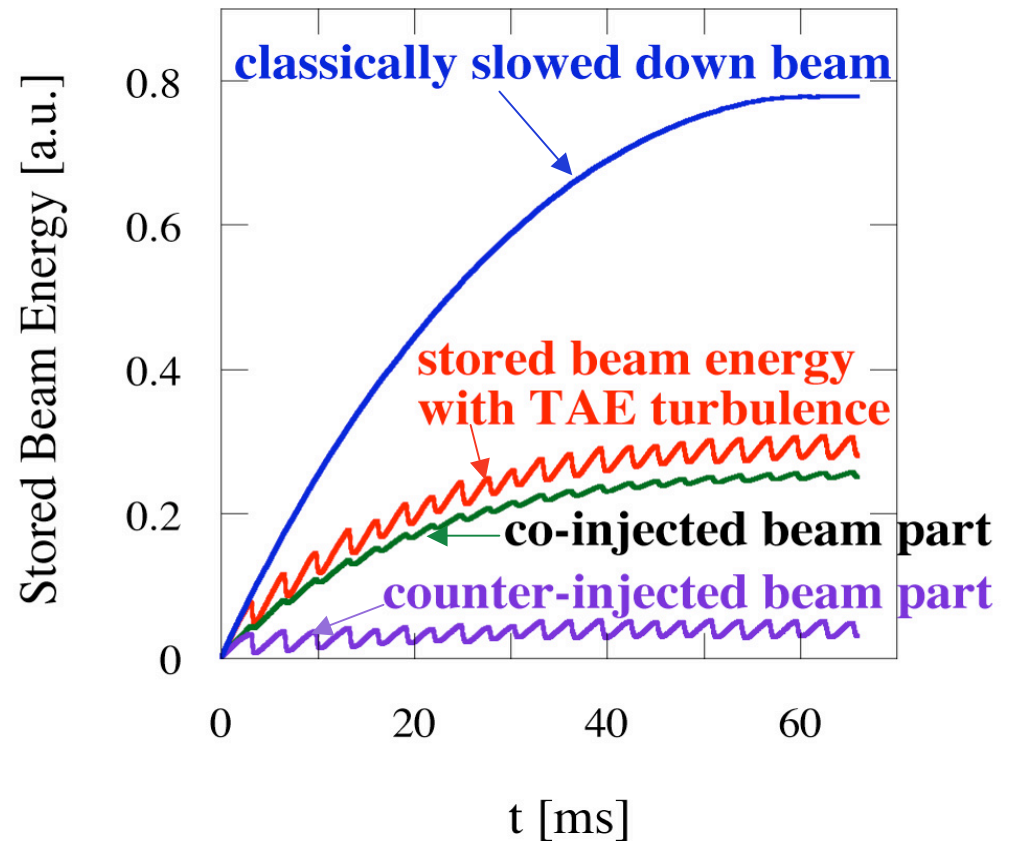
- Experiments show both benign and deleterious effects
- Rapid losses in early TAE experiments:
  - Wong (TFTR)
  - Heidbrink (DIII-D)
- Simulation of rapid loss:
  - Todo, Berk, and Breizman, Phys. Plasmas (2003)
  - Multiple modes ( $n=1, 2, 3$ )



K.L. Wong et al., PRL (1991)

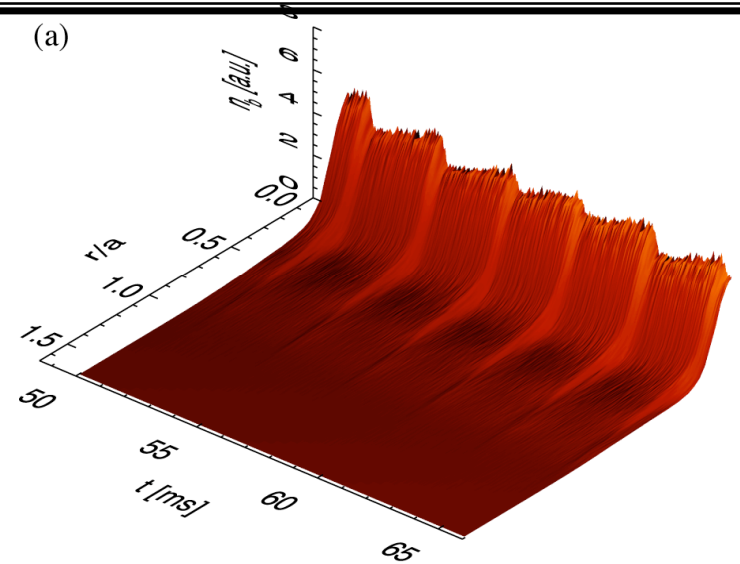
# SIMULATION OF INTERMITTENT LOSSES

- Simulations reproduce NBI beam ion loss in TFTR
- Synchronized TAE bursts:
  - At 2.9 ms time intervals (cf. 2.2 ms in experiments)
  - Beam energy 10% modulation per burst (cf. 7% in experiment)
- TAE activity reduces stored beam energy wrt to that for classical slowing-down ions
  - 40% for co-injected ions
  - Larger reduction (by 88%) for counter-injected beam ions (due to orbit position wrt limiter)

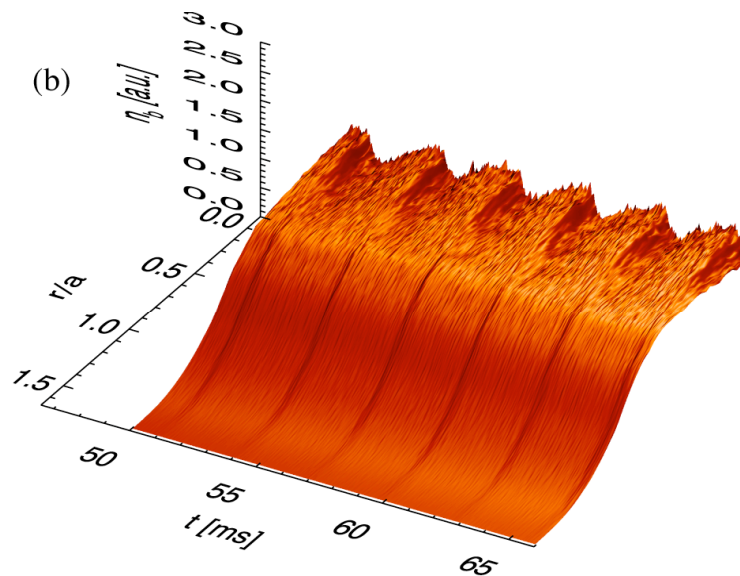


Y. Todo et al., Phys Plasmas (2003)

# TEMPORAL RELAXATION OF RADIAL PROFILE

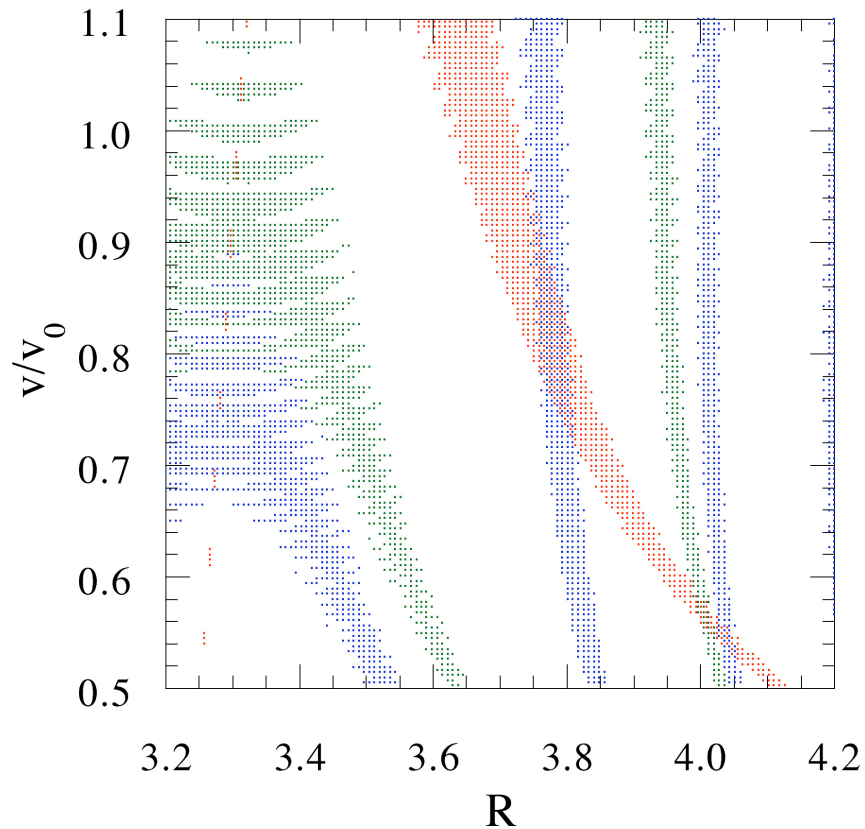


- **Counter-injected beam ions:**
  - Confined only near plasma axis



- **Co-injected beam ions:**
  - Well confined
  - Pressure gradient periodically collapses at criticality
  - Large pressure gradient is sustained toward plasma edge

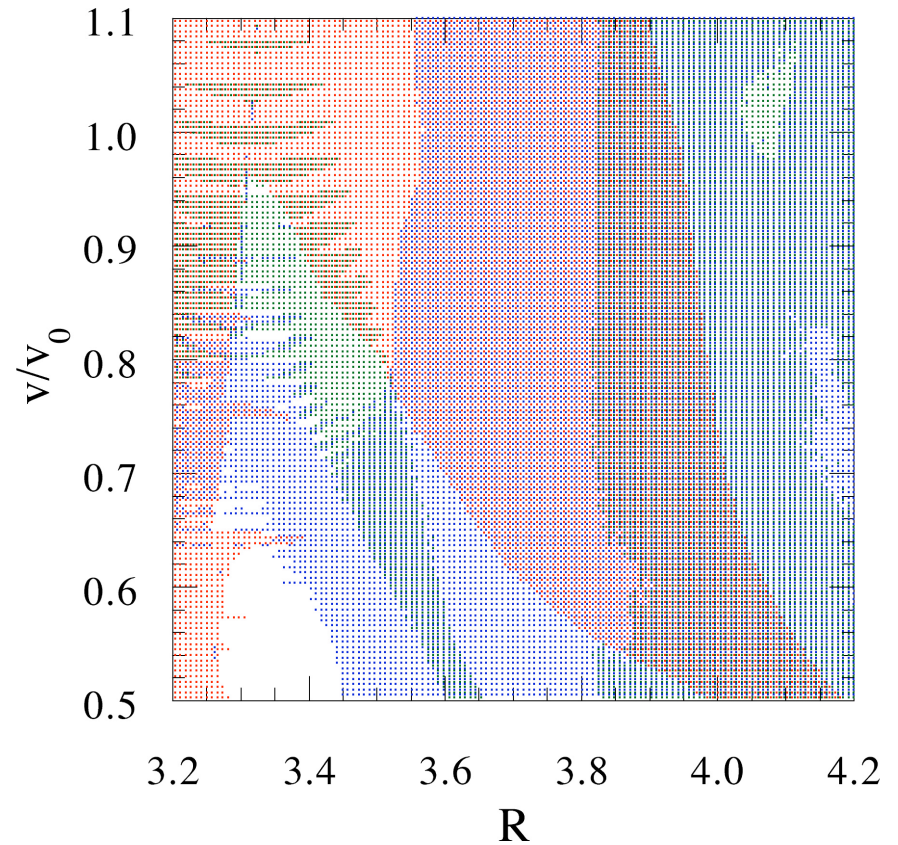
# PHASE SPACE RESONANCES



**For low amplitude modes:**

$$\square B/B = 1.5 \times 10^{-4}$$

$n=1$ ,  $n=2$ ,  $n=3$



**At mode saturation:**

$$\square B/B = 1.5 \times 10^{-2}$$

$n=1$ ,  $n=2$ ,  $n=3$

# ISSUES IN MODELING DIFFUSIVE TRANSPORT

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- **Reconciliation of mode saturation levels with experimental data**
  - Simulations (Y. Todo) reproduce experimental behavior for repetition rate and accumulation level
  - However, saturation amplitude is larger than exp'tal measurements
- **Edge effects in fast particle transport**
  - Sufficient to suppress modes locally near the edge
  - Need better description of edge plasma parameters
- **Transport barriers for marginally stable profiles**
- **Resonance overlap in 3D**
  - Different behavior: (1) strong beam anisotropy, (3) fewer resonances

# SUMMARY

# MODELING ISSUES FOR FURTHER STUDY

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- **Diagnostic applications of benign modes**
  - Obtain  $q_{\min}$  without saturation information
- **Quantitative interpretation of exp'tal data for nonlinearly saturated isolated modes (TAE, Cascade, high-frequency modes, etc.)**
  - What determines saturation level
- **Transport barriers for energetic particles in presence of many unstable modes**
- **Parameter window for fast particle confinement in fusion devices**
  - Predict range for safe operation
- **Intermittent strong bursts versus stochastic diffusive transport**
  - Many modes (quasilinear-like scenario)
- **Full lifecycle of fishbones and EPs, starting from instability threshold**
  - So far, described only the initial phase
- **Fast nonlinear frequency sweeping**
  - Multi-mode case
- **Interplay of kinetic and fluid resonances**

# MODELING/SIMULATION METHODOLOGY

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- **Highlight puzzles/discrepancies of fundamental importance**
- **Identify key physics ingredients qualitatively/analytically**
- **Develop self-consistent reduced theory-based models/codes:**
  - To capture and simulate essential physics
  - To perform broad parameter scans
- **Develop advanced techniques for large-scale simulations**
- **Use reduced models/codes to validate large codes**
- **Apply large codes:**
  - To quantitatively model present-day experiments
  - To predictively extrapolate to reactor conditions
- **Attain ultimate goal of integrated simulations**