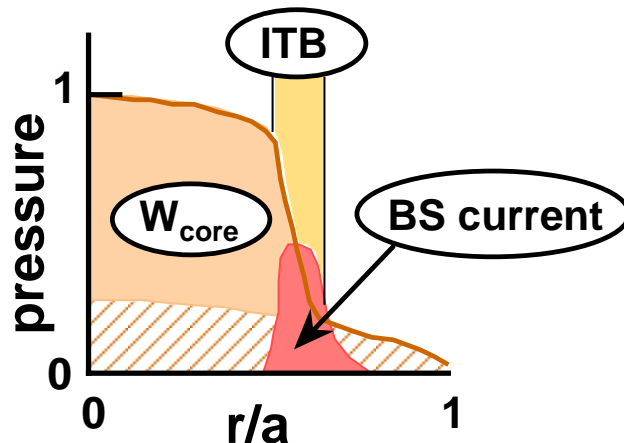


Axisymmetric Tri-Magnetic-Islands Equilibrium of Strongly-Reversed-Shear Tokamak Plasma - An Idea for the Current Hole -

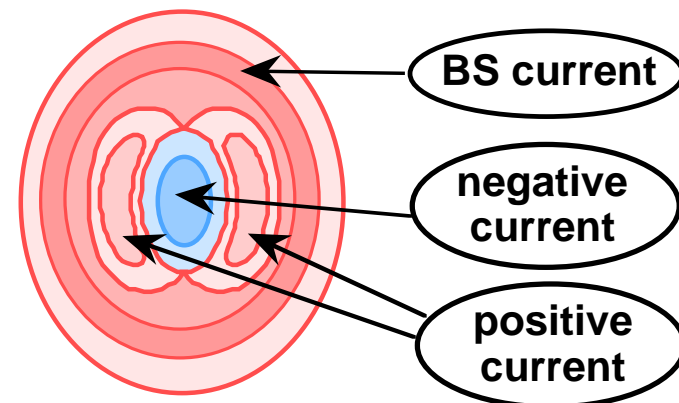
滝塚 知典 (原研那珂研)

Content

Confinement and Transport
of RS Plasmas with ITB

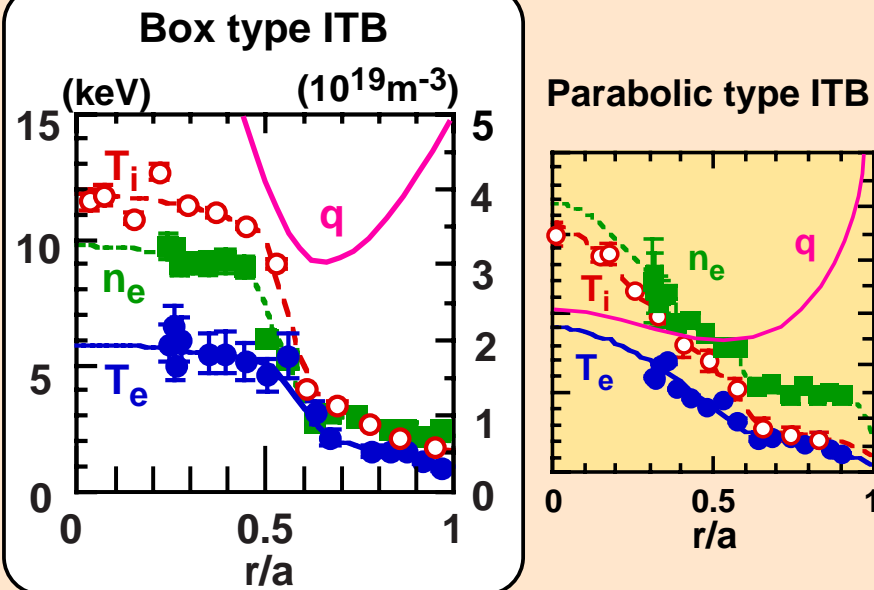


Axisymmetric Tri-Magnetic-Islands
(ATMI) Equilibrium

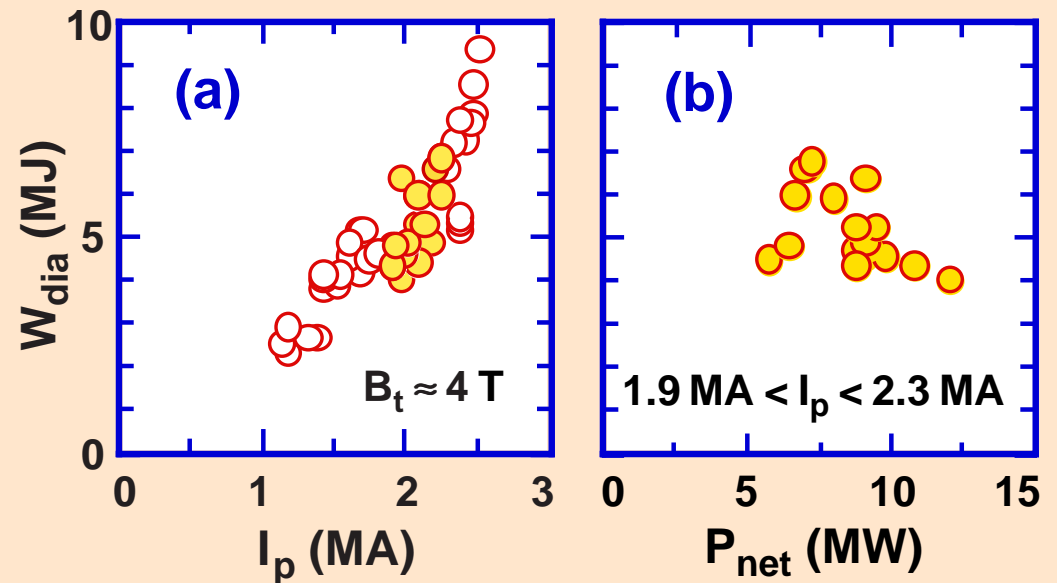


Confinement and Transport of RS Plasmas with ITB

T. Takizuka et al., "Energy Confinement Scaling for Reversed-Shear Plasmas with Internal Transport Barrier in JT-60U", Plasma Phys. Control. Fusion 44 (2002) A423.



JT-60U database
L-mode edge, Box type ITB

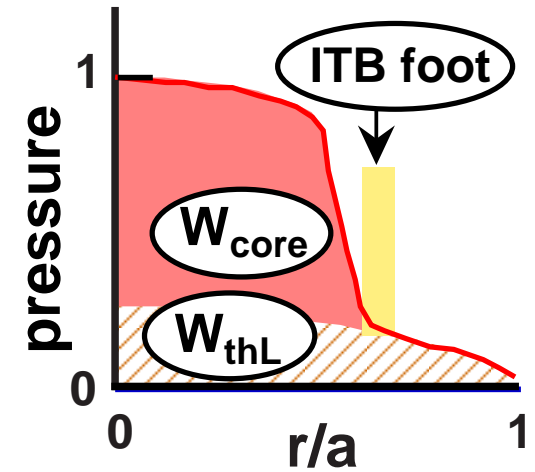


Stored energy depends strongly on plasma current I_p ($W_{\text{dia}} \sim I_p^{1.4}$), but little on heating power P_{net} ($\equiv P_{\text{NB}} - dW/dt$)

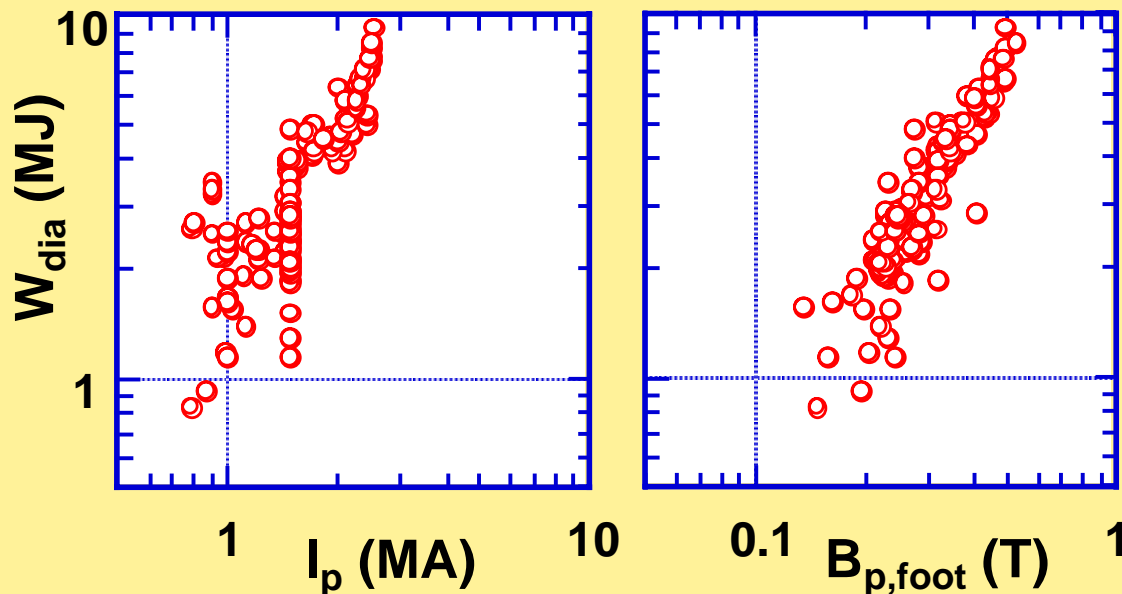
Working hypotheses

- (1) Stored energy is governed not by the heat flux through ITB but by the plasma parameters around ITB.
- (2) Stored energy is divided into L-mode base part and core part inside ITB.

$$[W_{\text{thL}} = 0.026 M^{0.3} I_p^{0.5} B_t^{0.39} R^{1.44} a^{0.92} \kappa^{0.88} n_{19}^{0.5} P^{0.33}]$$



W_{dia} depends more strongly on $B_{p,\text{foot}}$ than on I_p



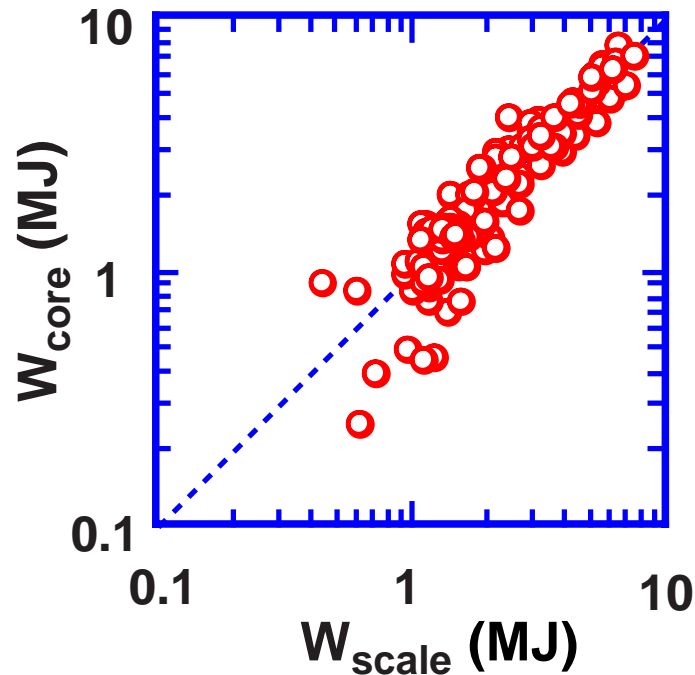
W_{core} can be expressed by a dimensionally-correct form without including P :

$$W_{\text{core}} \sim B_{p,f}^2 \times V_{\text{core}}$$

Energy Confinement Scaling

$$W_{\text{scale}} = C \epsilon_f^{-1} B_{p,f}^2 V_{\text{core}}$$

Analysis of Database
under the above hypotheses



$r = 0.95, \sigma = 22\%$

$$\epsilon_f \beta_{p,\text{core}} = \text{const. } (\approx 1/4)^*$$

Indication of MHD
Equilibrium Limit

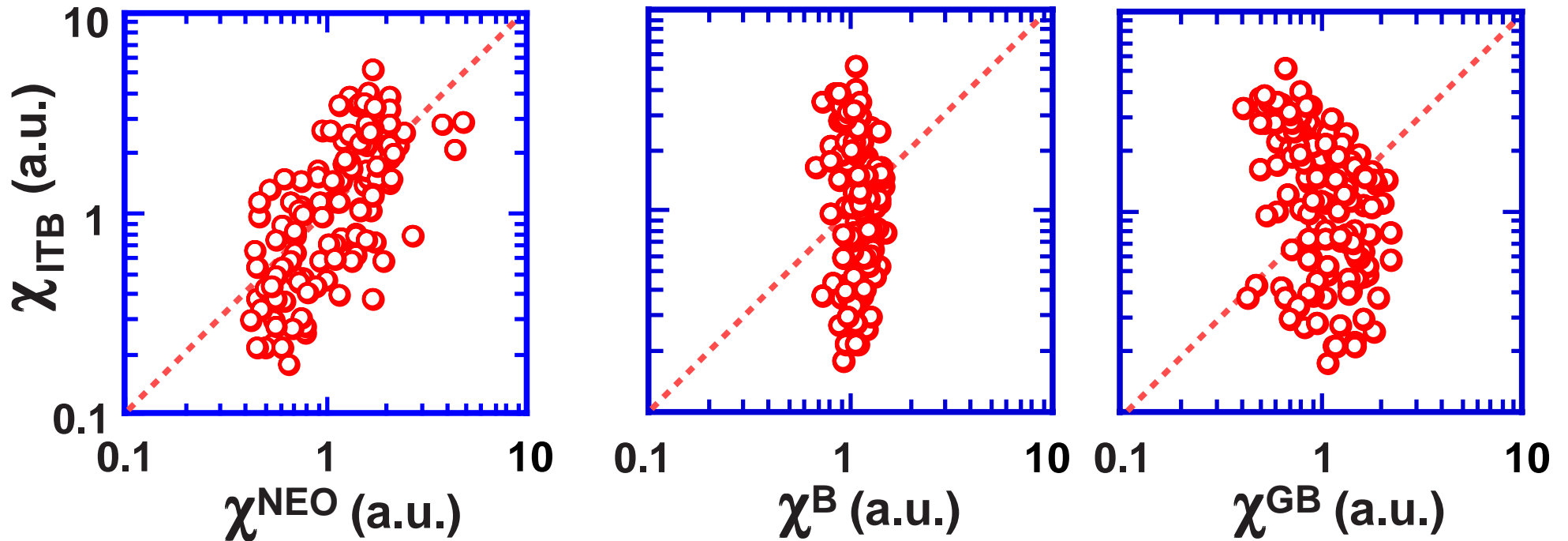
B_t : 2.1 ~ 4.1 T
 I_p : 0.8 ~ 2.6 MA
 P_{NB} : 4 ~ 19 MW
 ϵ_f : 0.12 ~ 0.19
 V_{core} : 9 ~ 25 m³
 W_{core} : 0.3 ~ 7.7 MJ

* $B_{p,f}$ is measured at
the outer midplane.
When averaged $B_{p,f}$
used,

$$\epsilon_f \beta_{p,\text{core}} > 1/2 .$$

"Neoclassical transport of ITB" is not inconsistent with the data, though the core stored energy is independent of P_{net}

$$\chi_{\text{ITB}} \propto P_{\text{net}} / (S_f n_{e0} |dT/dr|_{\text{ITB}}) \quad \chi^{\text{NEO}} \propto \epsilon_f^{0.5} n_{e0} / B_{p,f}^2 T_{i0}^{0.5}$$

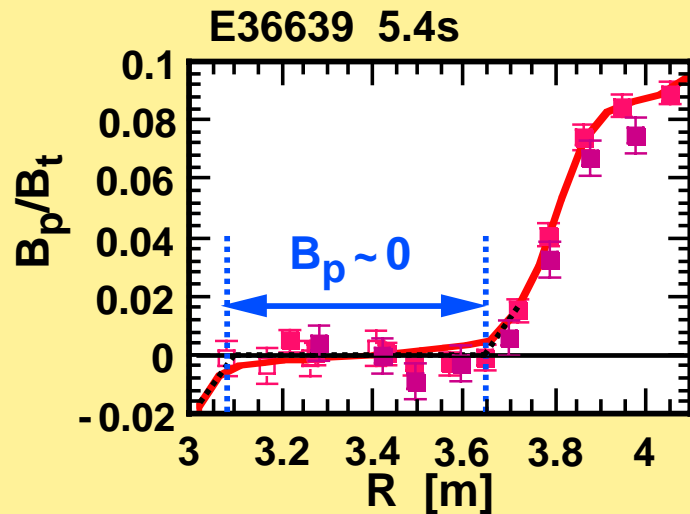


Bohm or GyroBohm transport is inconsistent with the data

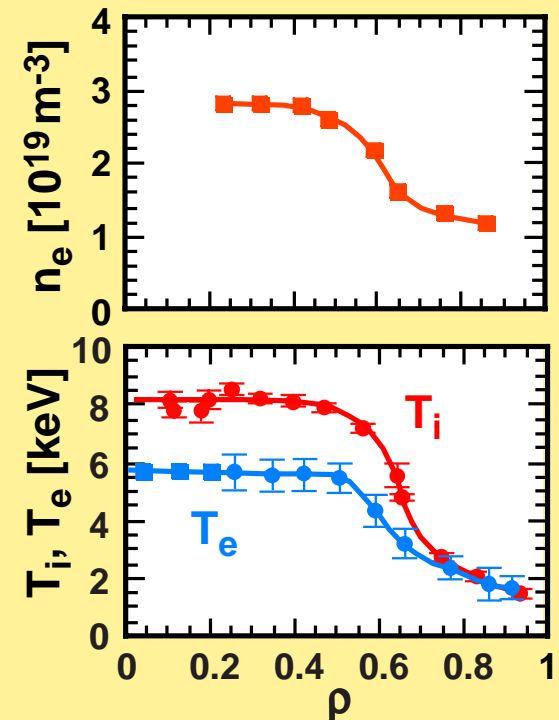
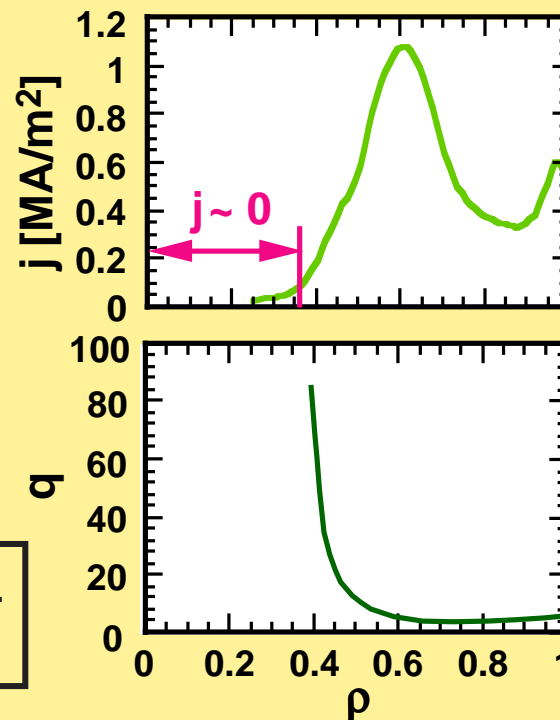
Axisymmetric Tri-Magnetic-Islands (ATMI) Equilibrium - An Idea for the Current Hole -

T. Takizuka, "Axisymmetric Tri-Magnetic-Islands Equilibrium of Strongly-Reversed-Shear Tokamak Plasma: An Idea for the Current Hole", J. Plasma Fusion. Res. 78 (2002) 1282;
<http://jspf.nifs.ac.jp/RCPDF/>

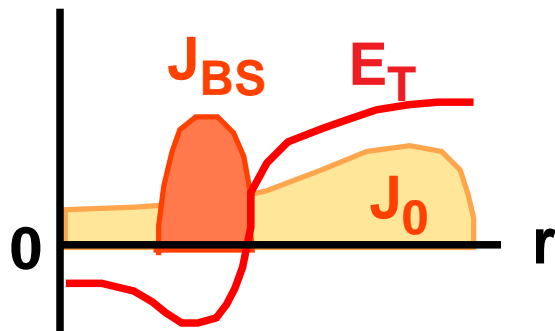
Current Hole in JT-60U Tokamak Plasma



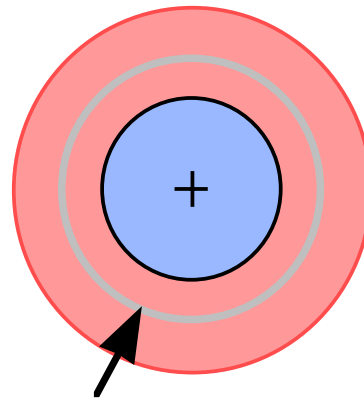
T. Fujita et al., Phys. Rev. Lett.
87 (2001) 245001



Generation of Negative Current and Instability of a Single-Magnetic-Axis (SMA) Equilibrium with Central Negative Current



Growth of bootstrap current J_{BS} induces negative E_T at the central region.

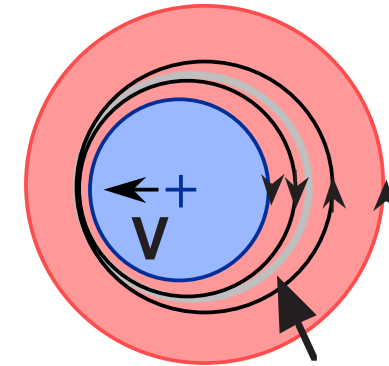


$$B_p = 0 ; \beta_p, I_i \rightarrow \infty$$

Central negative current in a SMA configuration destroys the equilibrium.

Large Shafranov shift

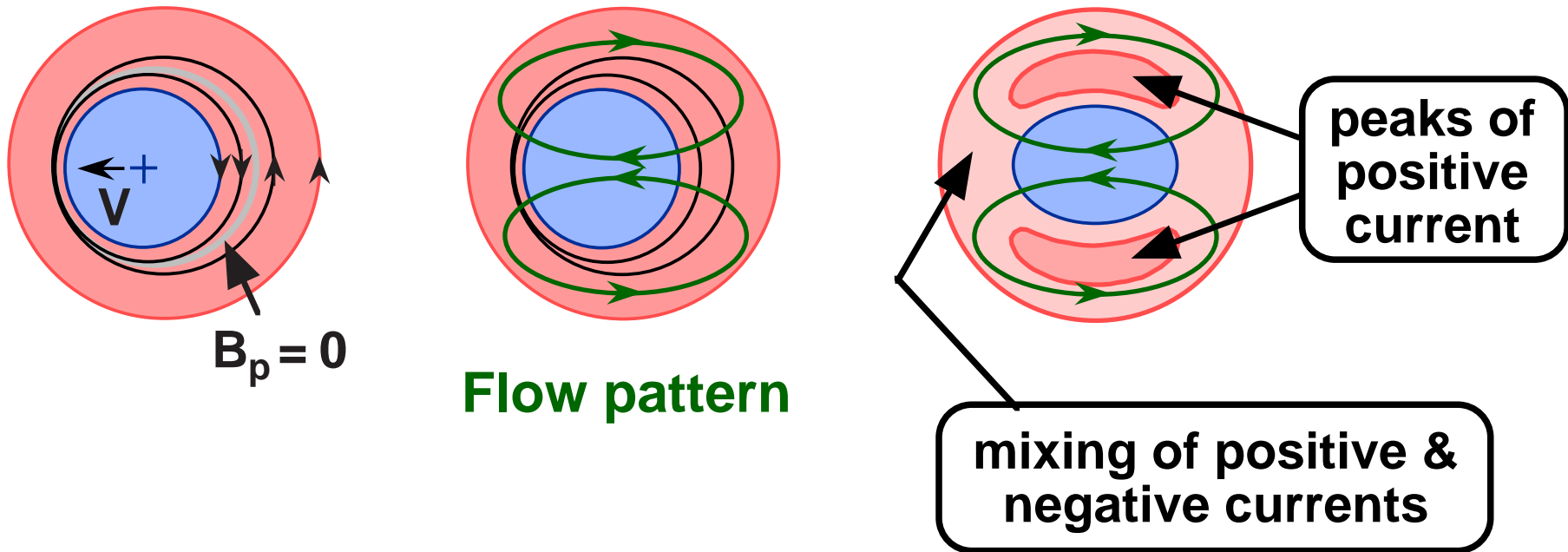
$$\Delta \sim \varepsilon \Lambda \quad (\Lambda = \beta_p + I_i/2 - 1)$$



$$B_p = 0$$

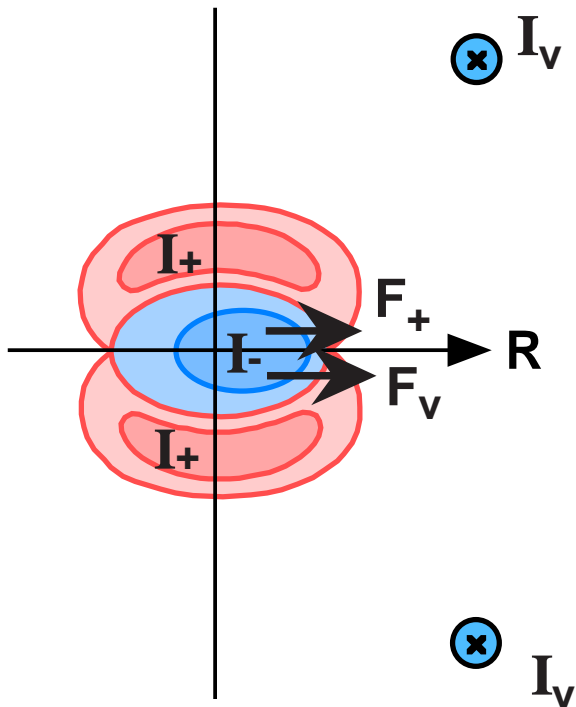
Even in a cylindrical plasma, central negative current in a SMA configuration destabilizes $n=0/m=1$ resistive MHD mode.

**Formation of Three Magnetic Islands ($m=2/n=0$)
through Nonlinear Evolution of $m=1/n=0$ Instability (?)**

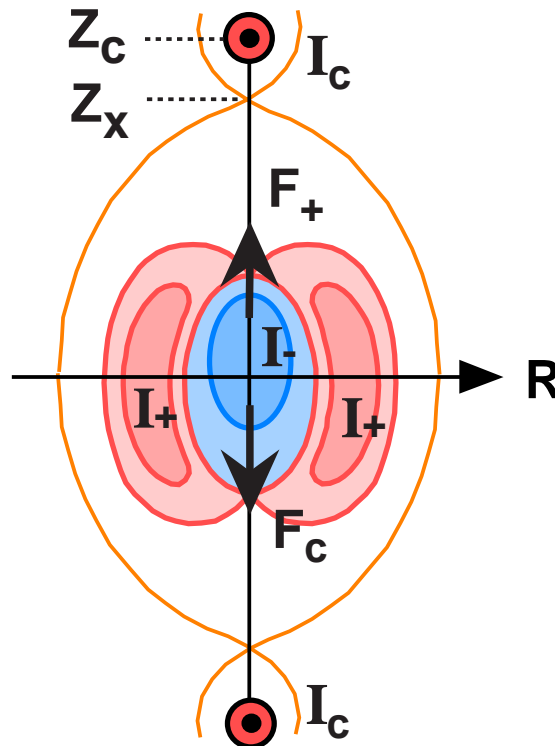


Unstable Configuration (Z arrangement) and Stable Configuration (R arrangement) under the Current Limit

Unstable Z-ar.



Stable R-ar.



Stability Condition (current limit)

$$F_+ < F_c$$

$$I_+ / (r_1^2 + \delta Z^2) < I_c / Z_c^2$$

$$I_+ / r_1^2 < I_p (Z_c - Z_x) / Z_x Z_c^2$$

$$q_{ATMI} / q_a > \kappa_1^2 Z_c^2 / (Z_c - Z_x) a$$

$$2I_+ / I_- \approx (1 + \kappa_1^2) / \kappa_1^2 Z$$

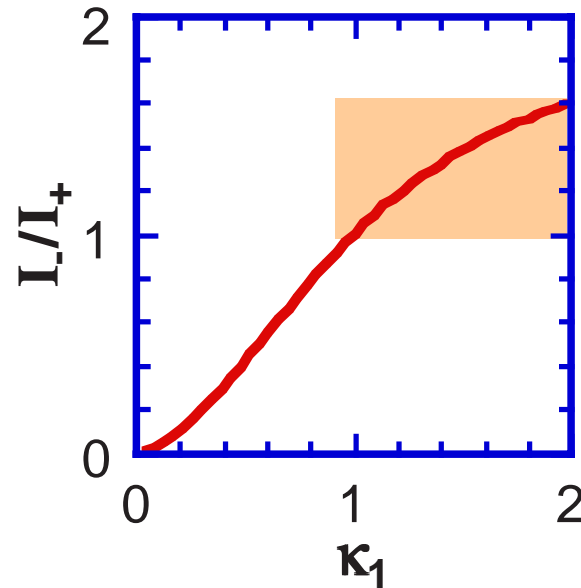
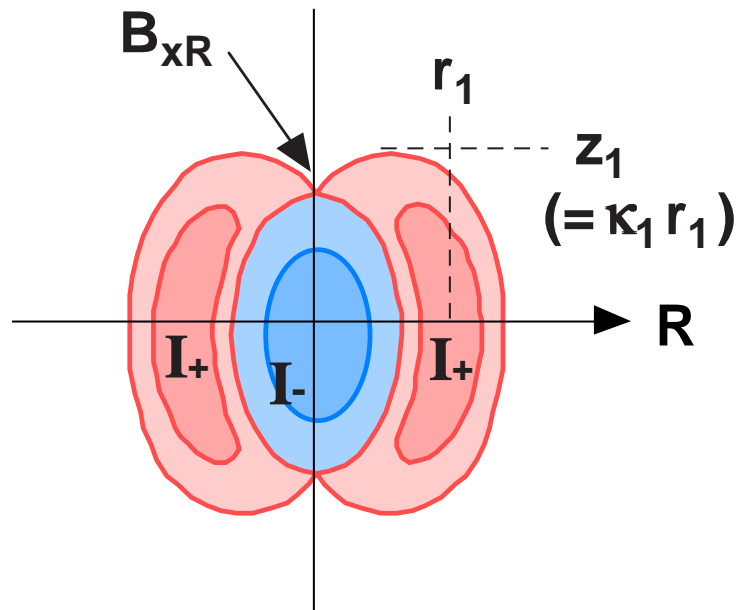
(see 2 page later)

Scenario for the Rotation from Unstable Z-ar. Configuration to Stable R-ar. Configuration

Rotation is not simply explained by MHD theory.

1. Very small B_p at the central region
2. Large ion-orbit size
3. Large radial electric field E
4. $E \times B$ rotation from Z-ar. to R-ar.
5. Closed magnetic surfaces of tri-magnetic islands at stable R-ar.
6. Uniform electrostatic potential on each magnetic surface
7. $E \times B$ rotation is on the surface, and stable R-ar. does not change its magnetic configuration.

Ratio of negative to positive currents in ATMI $I_-/I_+ \approx 1$



$$B_{xR} \propto 2I_+z_1/(r_1^2 + z_1^2) - I_-/z_1$$

$$B_{xR} = 0$$

$$I_-/I_+ \approx 2\kappa_1^2/(1 + \kappa_1^2)$$

Unbalance current flows out through x-point region by magnetic reconnection process.

Current drive for both co- and ctr- directions is difficult in the current hole region. (ECCD experiment in JT-60U)

Summary

Core Stored Energy : MHD equilibrium limit $\epsilon_f \beta_{p,core} = \text{const.}$

Neoclassical Transport at ITB

Axisymmetric Tri-Magnetic-Islands (ATMI) Equilibrium

Generation of central negative E_T due to the growth of J_{BS}

Formation of ATMI configuration through the nonlinear evolution of $m=1/n=0$ MHD instability (?)

ExB rotation to the stable R-ar. configuration (?)

Limit of the +/- current ratio $I_-/I_+ \sim 1$

Difficulty of current drive in the ATMI region

Limit of the ATMI current $q_{ATMI}/q_a > \kappa_1^2 Z_c^2 / (Z_c - Z_x) a$

$q_{ATMI} > 40$ for the JT-60U parameters

Current Hole is ATMI Equilibrium

Plasma profile in central region is not flat, but wavy along R direction